Euphotic zone depth: Its derivation and implication to ocean-color remote sensing

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Received 6 July 2006; revised 12 October 2006; accepted 1 November 2006; published 16 March 2007.

Euphotic zone depth, \( z_{1\%} \), reflects the depth where photosynthetic available radiation (PAR) is 1% of its surface value. The value of \( z_{1\%} \) is a measure of water clarity, which is an important parameter regarding ecosystems. Based on the Case-1 water assumption, \( z_{1\%} \) can be estimated empirically from the remotely derived concentration of chlorophyll-a ([Chl]), commonly retrieved by employing band ratios of remote sensing reflectance (\( R_{rs} \)).

Recently, a model based on water’s inherent optical properties (IOPs) has been developed to describe the vertical attenuation of visible solar radiation. Since IOPs can be near-analytically calculated from \( R_{rs} \), so too can \( z_{1\%} \). In this study, for measurements made over three different regions and at different seasons (\( z_{1\%} \) were in a range of 4.3–82.0 m with [Chl] ranging from 0.07 to 49.4 mg/m³), \( z_{1\%} \) calculated from \( R_{rs} \) was compared with \( z_{1\%} \) from in situ measured PAR profiles. It is found that the \( z_{1\%} \) values calculated via \( R_{rs} \)-derived IOPs are, on average, within ~14% of the measured values, and similar results were obtained for depths of 10% and 50% of surface PAR. In comparison, however, the error was ~33% when \( z_{1\%} \) is calculated via \( R_{rs} \)-derived [Chl]. Further, the importance of deriving euphotic zone depth from satellite ocean-color remote sensing is discussed.


1. Introduction

Euphotic zone depth, \( z_{1\%} \), reflects the depth where only 1% of the surface photosynthetic available radiation (PAR) remains [Kirk, 1994]. \( z_{1\%} \) is a measure of water clarity, which is not only a quality index of an ecosystem but also an important property for primary production [Behrenfeld and Falkowski, 1997; Platt and Sathyendranath, 1988] and heat transfer [Chang and Dickey, 2004; Kara et al., 2005; Sathyendranath et al., 1991] in the upper water column. Based on Case-1 water assumptions [Morel, 1988; Morel and Prieur, 1977], \( z_{1\%} \) can be estimated from remotely derived concentration of chlorophyll-a or total pigments, with either a spectrally integrated form [Morel, 1988] or a spectrally resolved formalism [Prieur and Sathyendranath, 1981; Sathyendranath et al., 1989a].

Case-1 waters are those whose inherent optical properties [Preisendorfer, 1976] can be adequately described by phytoplankton (represented by chlorophyll concentration) [Gordon and Morel, 1983; IOCCG, 2000; Morel, 1988; Morel and Prieur, 1977]. Case-1 waters thus require the optical properties of other optically active constituents (such as Colored Dissolved Organic Matter and suspended particles) closely follow the optical properties of phytoplankton [Morel, 1988; Morel and Maritorena, 2001]. Because the definition of Case-1 water is not based on the geographical location, nor based on the value of chlorophyll, it is difficult to know a priori if a water body fits the Case-1 definition when it is measured by a remote sensor. Frequently open ocean waters are assumed as Case-1, but recent studies [Lee and Hu, 2006; Mobley et al., 2004] have shown that this could be problematic, as in general water’s optical properties are not determined by phytoplankton alone [Sathyendranath et al., 1989b].

To incorporate the dependence of subsurface light field on other water constituents (such as CDOM), Liu et al. [2002] developed a numerical model (via look-up-table) that uses information about chlorophyll concentration, CDOM absorption and particle scattering coefficient as inputs to describe the vertical distribution of downwelling irradiance. This approach, similar as those based on Case-1 assumption [Morel and Antoine, 1994; Ohlmann and Siegel, 2000; Sathyendranath et al., 1989a], however, still requires accurate information about chlorophyll concentration when applied to ocean-color remote sensing. Liu et al. [2006] are successful using this approach for waters of the West Florida Shelf.

In another empirical approach, Mueller and Lange [1989] developed relationships to calculate \( z_{1\%} \) directly from the diffuse attenuation coefficient at 490 nm. The associated empirical constants, however, were developed...
for a specific region and specific temporal ranges [Mueller and Lange, 1989], it is then difficult to apply the empirical relationships to the global oceans at different seasons.

Recently, based on the principle that the vertical variation of subsurface light field is determined by inherent optical properties (IOPs), Lee et al. [2005b] developed an analytical model to describe the vertical attenuation of downwelling vector irradiance in the visible domain ($E_{VIS}$, 350–700 nm). In this IOP-centered approach, the vertical attenuation coefficient of $E_{VIS}$ ($K_{VIS}(z)$) is modeled as an analytical function of water’s absorption and backscattering coefficients. For vertical distributions of $E_{VIS}(z)$ ranging from 0.1 to 100% of the surface value, the modeled $E_{VIS}(z)$ is accurate to within ~6% of the actual value [Lee et al., 2005b], for a data set simulated using Hydrolight [Mobley, 1995]. Since the vertical distribution of $E_{VIS}(z)$ is nearly identical to that of $PAR(z)$ [Lee et al., 2005b; Morel and Gentili, 2004], $z_{10\%}$ can be easily calculated when IOPs (the absorption and backscattering coefficients at 490 nm, in particular) are known, either from in situ measurements or from remote sensing of ocean color.

Regardless of whether one uses the Case-1 approaches or the recent IOP-centered approach, there has been no test or validation of $z_{10\%}$ algorithms with a broad-range in situ data. On the other hand, global or basin scale estimation of $z_{10\%}$ from ocean color has been incorporated into studies of new production [Behrenfeld et al., 2005] and the Sverdrup hypothesis [Siegel et al., 2002]. To ensure reliable products for these biogeochemical studies, it is desirable to know the accuracy of remotely derived $z_{10\%}$ for broad range of waters.

In this study, for measurements made in the Arabian Sea, the Monterey Bay and the Gulf of Mexico that covered both oceanic and coastal waters and measured at different seasons, $z_{1\%}$ (and $z_{10\%}$, $z_{50\%}$) are calculated from spectral remote-sensing reflectance using both the Case-1 and IOP-centered approaches. With an emphasis on computational efficiency, only spectrally integrated approaches are studied here. The calculated $z_{x\%}$ ($x$ represents 1, 10, or 50 in this article) values are then compared with those from profiles of PAR measurements to evaluate the performance of the approaches. Further, we discussed the unique characteristics of property $z_{1\%}$ (and/or $z_{10\%}$) and its importance in ocean color remote sensing.

2. Data and Methods

Field measurements from three different regions are used for testing and evaluating the methods for deriving $z_{x\%}$ from $R_{PAR}$. These regions are: the Monterey Bay (data collected in September and October 1989), the Gulf of Mexico (April and June 1993), and the Arabian Sea (December 1994). These water environments include open-ocean oligotrophic waters (the Loop Current, and Arabian Sea), coastal high-productivity waters (the Monterey Bay, the West Florida Shelf), and turbid Mississippi River plume waters. The range of chlorophyll-a concentration measured was ~0.07–49.4 mg/m$^3$ with $z_{1\%}$ ranging from ~4.3 to 82.0 m. During these measurements the above-surface solar zenith angles ($\theta_s$) were between 8° and 80° (30% of them with $\theta_s > 50°$). For the conditions with the Sun covered by clouds, $\theta_s$ is taken in our calculations as 45° as in Sathyendranath et al. [1989a].

There are 65 stations that include measurements of both remote-sensing reflectance and vertical profiles of PAR. Remote-sensing reflectance at wavelength $\lambda$, $R_{PAR}(\lambda)$, was calculated from measurements made above the sea surface as described in Carder and Steward [1985], with upwelling radiance, downwelling irradiance, and downwelling sky radiance measured by a handheld spectroradiometer. The methodology of determining in situ $R_{PAR}(\lambda)$ is described in detail in NASA protocols [Mueller et al., 2002a].

Instantaneous PAR (400–700 nm) in the upper water column at time $t$ ($PAR(z)$) was measured by lowering a PAR sensor ( Biospherical Instruments, Inc.) from surface to depth (z m, positive downward). $z$ was measured by a pressure sensor, and was accurate to within 0.1 m with a fall rate of 1 m per second. During the deployment of the PAR sensor in the water column, a deck cell was operated to measure simultaneously the above-surface downwelling irradiance at time $t$ ($E_s(490, t)$), in order to correct for any variations of input solar light resulting from passing clouds. The vertical profile of PAR free of cloud effects is [Smith et al., 1984]

$$PAR(z) = PAR_0(z) \frac{E_s(490 \text{ nm}, 0)}{E_s(490 \text{ nm}, t)}$$

From these $PAR(z)$ values, the ratio of PAR at depth to surface PAR $PAR(0)$ was calculated as

$$r_{PAR} = \frac{PAR(z)}{PAR(0)}$$

$z_{50\%}$, $z_{10\%}$, and $z_{1\%}$ (corresponding to $r_{PAR}$ equals 0.5, 0.1 and 0.01, respectively) were then determined from the vertical profile of $r_{PAR}(z)$. It is rare to have recorded depth with $r_{PAR}(z)$ exactly the desired ratio (10.0%, for example). The measured $z_{1\%}$ value is then an approximation by exponentially interpolating $r_{PAR}(z)$ between ~1.1x% and 0.9x%. Because the depth interval reported for the vertical profiler is less than 1 m, the maximum error in $z_{1\%}$ is less than 0.5 m.

For some earlier measurements (the 9 Monterey Bay stations in 1989 and 10 stations in the Gulf of Mexico in 1993), the PAR sensor was not sensitive enough to provide a reading when PAR is 1% of $PAR(0)$. For those stations, only $z_{10\%}$ and $z_{50\%}$ were directly obtained from $r_{PAR}(z)$ profiles.

For the other stations where both $z_{10\%}$ and $z_{1\%}$ could be obtained from $r_{PAR}(z)$, it is found that a good linear relationship exists between $z_{10\%}$ and $z_{1\%}$ (see Figure 1). For $z_{1\%}$ in the range of 4.3–64.4 m ($z_{1\%}$ in the range of 2.1–29.0 m), by forcing the linear-regression line to pass through the origin, the following relation was obtained,

$$z_{1\%} = 2.25 z_{10\%}$$

with a coefficient of determination of 0.99 ($n = 43$). In this regression analysis, the two measurements with $z_{1\%}$ of ~82.0 m (the black points in Figure 1) were excluded, simply because they were significantly beyond the general linear trend with too few points in that range to make a
confident generalization. For the same ranges of \(z_{1\%}\) and \(z_{10\%}\), simulations using Hydrolight [Mobley, 1995] with different sun angles and IOPs indicate a nearly identical relationship as equation (3) (with a slope of 2.21). These results then suggest that equation (3) does represent a reliable relationship between \(z_{1\%}\) and \(z_{10\%}\), at least for \(z_{10\%}\) in this specific range of \(\sim 2\text{–}30\) m. Therefore, for those stations where no \(z_{1\%}\) were obtainable from \(r_{PAR}(z)\) and \(z_{10\%}\) was in the range of \(2\text{–}30\) m, \(z_{1\%}\) values representing PAR profiles were derived by applying equation (3). In essence, in this study, equation (3) is only utilized for range-constrained interpolation, not for extrapolation.

[14] Usually, \(z_{10\%}\) is assumed to be half of \(z_{10\%}\) (\(z_{10\%}\) is called midpoint depth in Kirk [1994]). This is based on the assumption that the attenuation coefficient for PAR (\(K_{PAR}\)) does not change with depth. However, because PAR converges to wavelengths with less attenuation coefficients when PAR propagates from surface to deeper depths, \(K_{PAR}\) is always smaller at depth than at surface (the difference can be a factor of 2 or 3) even for a vertically homogeneous water environment [Lee et al., 2005b; Morel, 1988; Zaneveld et al., 1993], and thus \(z_{10\%}\) is shallower than half of \(z_{1\%}\).

3. Derive \(z_{1\%}\) From \(R_{rs}\)

3.1. IOP-Centered Approach

[15] The vertical distribution of downwelling irradiance in the visible domain \((E_{VIS}, 350\text{–}700\) nm) is described as

\[
E_{VIS}(z) = E_{VIS}(0)e^{-K_{VIS}(z)}. \tag{4}
\]

\(K_{VIS}(z)\), the attenuation coefficient of \(E_{VIS}\), however, is no longer treated as independent of \(z\) but a function of both \(z\) and inherent optical properties [Lee et al., 2005b],

\[
K_{VIS}(z) = K_1 + \frac{K_2}{(1 + z)^b}, \tag{5}
\]

with

\[
\begin{align*}
K_1 &= \left[\chi_0 + \chi_1(a(490))^{0.5} + \chi_2 b_0(490)\right](1 + \alpha_0 \sin(\theta_s)), \\
K_2 &= [\zeta_0 + \zeta_1 a(490) + \zeta_2 b_0(490)](\alpha_1 + \alpha_0 \cos(\theta_s)),
\end{align*}
\]

where \(a(490)\) and \(b_0(490)\) are water absorption and backscattering coefficients at 490 nm, and \(\chi_0, \chi_1, \chi_2, \zeta_0, \zeta_1, \zeta_2\) and \(\alpha_0, \alpha_1, \alpha_2\) are model constants derived from Hydrolight simulations [Lee et al., 2005b]. Given a sun angle and values of \(a(490)\) and \(b_0(490)\), \(K_{VIS}(z)\) can then be calculated for any depth.

[16] Defining the optical depth, \(\tau_E\), for \(E_{VIS}\) as

\[
K_{VIS}(z)z = \tau_E, \tag{7}
\]

then the depth where \(E_{VIS}(z) = 1\%\) of \(E_{VIS}(0)\) is \(-\ln(0.01) = \tau_E = 4.605\). Combining equations (5) and (7), after simple math manipulations, a cubic-polynomial equation with \(z\) as the variable is obtained

\[
z^3 + y_1 z^2 + y_2 z + y_3 = 0, \tag{8}
\]

with \(y_1 \text{–} 3\) functions of \(K_1\), \(K_2\) and \(\tau_E\),

\[
\begin{align*}
y_1 &= \frac{k_1^2 - k_2^2 - 2\tau_E k_1}{k_2^2}, \\
y_2 &= \frac{k_1^2 - 2\tau_E k_1}{k_2^2}, \\
y_3 &= \frac{k_1^2}{k_2^2}.
\end{align*}
\]

Mathematically there are three solutions (one negative, two positives) that satisfy equation (8), but it is the smaller, positive, value that is consistent with radiative transfer theory. In this study, as examples to demonstrate this IOP-centered approach, \(z\) for \(\tau_E\) equals 4.605 (\(E_{VIS}(z)/E_{VIS}(0) = 0.01\)).
1\%), 2.303 (10\%), and 0.693 (50\%) are calculated for those measurements, respectively.

[17] Before the derivation of ε for different τz, inherent optical properties (a(490) and b0(490) in particular) were derived from Rs by using the quasi-analytical algorithm (QAA) [Lee et al., 2002]. QAA is an algorithm for deriving absorption and backscattering coefficients of optically deep waters from ocean-color remote sensing. Its concept and architecture are documented in detail in Lee et al. [2002]. Here a slightly updated version of QAA (v4) was applied, and the details of the updates are provided in Appendix A.

[18] It is necessary to point out that EVIS in equation (4) represents downwelling vector irradiance in the range of 350–700 nm, and is measured by energy (W/m²), whereas the PAR sensor measures both downwelling and upwelling scalar irradiance in the range of 400–700 nm and is measured by the amount of photons (quanta/m²/s). For the same wavelength range (400–700 nm), Hydrolight simulations [Morel and Gentili, 2004] indicate that z1% is about the same when it is measured either by EVIS or by PAR (Table 2 of Morel and Gentili [2004]). Also, sensitivity tests with Hydrolight simulations indicate that the vertical profile of EVIS is nearly identical to that of PAR, because the upwelling irradiance is in the order of ~5% of downwelling irradiance, and that the irradiance in the range of 350–400 nm is small. In general, the difference between z1% measured by either EVIS or PAR is small (less than 10%). Therefore, z1% calculated from equation (4) is considered equivalent and comparable to that determined from the vertical profiles of PAR(z).

3.2. [Chl]-Centered Approach

[19] Based on the Case-1 water assumption, and in a spectrally integrated form, z1% can be estimated from chlorophyll-a concentration ([Chl]) of the surface layer (A. Morel, personal communication, after the statistical analysis of the shape of [Chl] vertical profiles [Morel and Berthon, 1989]),

\[ z_{1\%} = 34.0([\text{Chl}])^{-0.39}. \]  

(10)

Because surface [Chl] can be calculated from Rs(λ) by the operational OC4v4 algorithm [O’Reilly et al., 2000],

\[ [\text{Chl}] = 10^{0.366 - 3.067 \rho + 1.93 \rho^2 + 0.649 \rho^3 - 1.532 \rho^4}, \]  

(11)

with

\[ \rho = \log \left( \frac{\max(Rs(440, 490, 510))}{Rs(555)} \right), \]  

(12)

it is then quite straightforward to estimate z1% via this route.

4. Results and Discussion

[20] To evaluate and analyze the performance of deriving z1% from ocean-color remote sensing. Figure 2 (a for rPAR = 0.5 (50%), 2b for rPAR = 10%, and 2c for rPAR = 1%) and Figure 3 present Rs-derived values (εPAR) versus those determined from rPAR(z) (εPAR), and Table 1 summarizes their errors. Quantitatively, an averaged percentage error for all observations is calculated from

\[ \varepsilon = \frac{\sum_{i=1}^{n} |e_i - m_i|}{n} \times 100\%. \]  

(13)

[21] Measurements at one station were excluded from these analyses because z1% from remote-sensing reflectance (no matter which method) is about four times the z1% from PAR profile. This one station is regarded as an outlier because of the significant difference between in-water and above-water determinations.

4.1. IOP-Centered Approach

[22] For this data set (64 stations) that covers both oceanic and coastal waters, the average percentage error (ε) between modeled and measured property is 18.5% (with a maximum error of 70.9%) for z50% (for a range of 0.6–18.4 m from PAR measurements); while the errors are 13.8% (maximum error of 61.2%) for z10% (2.1–47.1 m), and 13.7% (maximum error of 63.5%) for z1% (4.3 to 82.0 m). The root-mean-square error in log scale (RMSE) is 0.079 for z1%, significantly smaller than the RMSE (0.329) of Rs-derived [Chl] (also see Figure 4).

[23] For z50%, excluding the residual errors in the KVIS model and the QAA algorithm, there are a few extra sources to contribute to its discrepancies. These extra sources include (1) PAR(z) attenuates sharply at surface [Paulson and Simpson, 1977; Zaneveld and Spinrad, 1980], making it harder to precisely determine z50% from the profile of rPAR(z); (2) surface layer suffers more influence from the wavy surface [Zaneveld et al., 2001], which can converge or diverge incoming solar radiation; and (3) variable ship-perturbation to the near-surface sensor [Gordon, 1985].

[24] Additionally, larger error came from the two clear-water stations (the two with circles in Figure 2), where surface chlorophyll-a values were around 0.07 mg/m³. The two stations had similar [Chl] values and similar Rs spectra, and nearly identical z10% and z1% depths from PAR(z) profiles (see Figures 2b and 2c); their z50% depths from PAR(z), however, differed by nearly 80%. It is not clear yet what might have caused such big differences in the z50% from PAR(z) measurements. Excluding these two stations, the average error for z50% dropped to 16.8%.

[25] Much better results are achieved for z10% (and z1%), though, again, larger errors occurred at the two clearer-water stations (Figure 2b, in circle). Unlike z50%, however, the z10% values from PAR(z) are about the same for these two stations. Overall, it appears that there is a trend of slight underestimation for z10% in the range of 10–30 m for this historical data set. High-quality measurements of both IOPs and PAR(z) profiles are required to isolate the error sources.

[26] For the entire range of z1% (4.3–82.0 m) measured at different times and from different regions, the Rs derived values clearly matched the PAR(z)-determined values very well. Though many potential sources of error prevent exact agreement between the two independent data sets, the low average error (13.7% in linear scale, 0.079 in RMSE) for z1% indirectly validated the approach of deriving IOPs from Rs and calculating KVIS(z) from IOPs. It is even more
encouraging given that the measurements were not made in vertically homogeneous waters. If we use the vertical variation (within the euphotic zone) of the diffuse attenuation coefficient [Kirk, 1994] at 440 nm ($K_d(440)$) as a proxy to represent the non-homogeneity of optical properties in the upper water column, the average coefficient of variation is 17% (maximum is 51%) for this data set. If the variation is determined by chlorophyll fluorescence profile, the aver-

Figure 2. Comparison between $z_{50\%}$ from IOP-centered approach and $z_{\text{PAR}}$ from in situ PAR profile (a) for $z_{50\%}$, (b) for $z_{10\%}$ and (c) for $z_{1\%}$. The two circled stations (Figure 2a) had nearly identical $z_{10\%}$ (Figure 2b) and $z_{1\%}$ (Figure 2c) values.
The age coefficient of variation is 23% (maximum is 91%). These values indicate varying optical or biological properties in the upper water column for each station. The excellent results of remotely estimated $z_{1\%}$ and $z_{10\%}$ apparently indicate that IOPs from the semianalytical algorithm and the analytical $K_{VIS}$ model integrate and largely compensate for some of the vertical variations. At least in part, this is because remote-sensing reflectance is always a measure of the weighted average of optical properties in the upper water column [Gordon and Clark, 1980; Sathyendranath and Platt, 1989; Zaneveld et al., 2005].

It is important to recognize that what was carried out here were not regressions from the data set to develop an empirical relationship, but rather were independent tests of a system developed earlier from other sources. Such efforts, together with the resultant small error, provide us much higher confidence in the application of such semianalytically based approaches to broader ranges of environments. Additionally, there is no requirement of the waters to be Case-1 for the application of this IOP-centered approach, thereby avoiding the hurdle of identifying a water body as Case-1 or not [Lee and Hu, 2006; Mobley et al., 2004]. Waters with shallower $z_{1\%}$ were often in coastal regions, and their optical properties were less likely to co-vary with chlorophyll concentration. This is also consistent with an earlier study regarding the downwelling diffuse attenuation coefficient [Lee et al., 2005a], where its value was underestimated (also based on $R_{rs}$-derived [Chl]) for higher values (shallower $z_{1\%}$), but was quite good for lower values (deeper $z_{1\%}$). Compared with the $z_{1\%}$ estimated from OC4v4-[Chl], $z_{1\%}$ estimated from the IOP-centered approach are generally shallower and more consistent with measurements, which would lead to a smaller compensation irradiance if derived as in Siegel et al. [2002].

It is necessary to emphasize that “the $z_{1\%}$ relationship developed in Morel [1988] and Morel and Maritorena [2001] requires either the mean chlorophyll concentration – or the water-column-integrated concentration – within the euphotic zone as input” (S. Maritorena, personal communication). Because of the existence of subsurface maxima of chlorophyll concentration, the mean value is normally greater than the surface value. Consequently, if surface chlorophyll (e.g., the product from current ocean-color

![Figure 3. Comparison between $z_{1\%}$ from [Chl]-centered approach and $z_{1\%}$ from in situ PAR profile. $z_{1\%}$ is calculated by equation (10), and all waters are assumed Case-1.](image)

<table>
<thead>
<tr>
<th>Table 1. Properties and Their Error Characters When Derived From Ocean-Color Remote Sensing</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>IOP-Centered Approach</strong></td>
</tr>
<tr>
<td>$z_{50%}$</td>
</tr>
<tr>
<td>$\varepsilon$ (average error)</td>
</tr>
<tr>
<td>Max. error</td>
</tr>
<tr>
<td>RMSE</td>
</tr>
</tbody>
</table>
remote-sensing algorithms) is used for the calculation of $z_{1\%}$, it is very likely that significant overestimation of $z_{1\%}$ values will result, as shown by Figure 5. Here $z_{1\%}$ is calculated simply using the relationship developed in Morel and Maritorena [2001]. Since we do not always know the details of the vertical distribution of chlorophyll concentration for each station (especially from remote sensing), it is assumed arbitrarily that the mean concentration of chlorophyll within the euphotic zone is 1.3 times that derived by the OC4v4 algorithm. With this consideration, the average error for $z_{1\%}$ is 34.3%. The average error is much larger ($\varepsilon = 49.9\%$) when no such adjustment is considered, but can be reduced to 22.3% if the mean concentration is assumed as 1.8 times the OC4v4 derived surface value.

5. Conclusions and Implication to Ocean-Color Remote Sensing

[30] In this study, euphotic zone depths (and two other environmental-optics depths) were derived semianalytically from spectral remote-sensing reflectance for measurements made in the Arabian Sea, the Gulf of Mexico, and the Monterey Bay. The derived values were then compared with the values from in-water vertical profiles of PAR. It was

**Figure 4.** Comparison between OC4v4 derived [Chl] and measured [Chl] from surface samples.

**Figure 5.** Comparison between $z_{1\%}$ from [Chl]-centered approach and $z_{1\%}$ from in situ PAR profile. Here $z_{1\%}$ is calculated by rearranging equation (6) of Morel and Maritorena [2001], and the average [Chl] in the euphotic zone is assumed as 1.3 times the surface value derived by OC4v4 (see text for details).
found that the two agreed with each other very well for the entire range of the data set. The average percentage error in the linear scale is 18.5% for \( z_{50\%} \) (in a range of 0.6–18.0 m), 13.8% for \( z_{10\%} \) (2.1–47.1 m), and 13.7% for \( z_{1\%} \) (4.3–82.0 m). Such small errors suggest a closure between the two independent measurements and determinations, and indirectly validate the semianalytical derivation of IOPs from \( R_s \) and the IOP-centered model of \( K_{\text{VIS}} \). To ensure its reliable applications to broad range of waters, however, more tests and validations with a wider dynamic range are certainly desired, especially for waters with \( z_{1\%} \) deeper than 100 m.

[31] Also presented are comparisons of \( z_{1\%} \) calculated from the same \( R_s \) but based on a methodology that uses chlorophyll values and the Case-1 water assumption. The average error is much larger (32.7%) when compared with \( z_{1\%} \) from measurements, with better results for oceanic waters, but overestimates \( z_{1\%} \) for coastal waters. As discussed earlier, the earlier source contributing to this larger error includes the inclusion of coastal waters with many stations hardly belonging to the Case-1 category. An average error of 75% in [Chl] determined by the OC4v4 algorithm for these waters suggests that use of a global algorithm for coastal waters imparts a significant contribution to the error field.

[32] The IOP-centered approach, however, worked not only with oceanic waters, but also with the more complex coastal and shelf waters, as represented by the data collected in this study. Also, because the relationships between \( z_{0\%} \) and \( R_s \) are linked analytically by IOPs, this IOP-centered approach avoids parameterizations regarding the wide regional and seasonal bio-optical variations. This is extremely important if an algorithm is going to be applied to get quantitative global observations. In the current process of deriving [Chl] of global oceans from satellite ocean-color remote sensing, the “global” algorithm is with a single set of parameters that does not correct for regional and seasonal bio-optical variations when applied to global oceans. If proper regional and seasonal parameterizations are utilized as desired [Carder et al., 1989; O’Reilly et al., 1998], a different global mean value for chlorophyll, and, likely, a different temporal trend of this mean relative to those of Gregg and Conkright [2002] and Antoine et al. [2005], would be reached (Shang and Lee, manuscript in preparation, 2007).

[33] Euphotic zone depth (\( z_{1\%} \)) (or 10%–light depth, \( z_{10\%} \)) represents depths where only 1% (or 10%) of surface PAR remains. Compared with the measurement of chlorophyll concentration (large uncertainties exist between measured by HPLC method and measured by fluorometric method [Mueller et al., 2002b; Trees et al., 1985]), or measurement of absorption and scattering coefficients of the bulk water or phytoplankton, \( z_{1\%} \) (or \( z_{10\%} \)) is much easier and more accurate to determine in the field. It does not require an absolute radiometric calibration of the PAR sensor, and it does not require delicate calibration of the final product or complicated and commonly error-bearing post processing (e.g., the absorption coefficient of phytoplankton from filter-pad technique [Allali et al., 1995; Cleveland and Weidemann, 1993; Mitchell, 1990]).

[34] Also, the close agreement between \( R_s \)-calculated and field-measured \( z_{1\%} \) (and \( z_{10\%} \)) can, in large part, be attributed to the fact that both are cumulative measurements of the upper water column. Chlorophyll concentrations or absorption and backscattering coefficients from discrete water samples may not accurately represent the average of the upper water column unless it is uniform; and, consequently, larger uncertainty will be introduced when comparing the values from discrete water samples with those from remote sensing. Many of those uncertainties would be minimized for \( z_{1\%} \) (or \( z_{10\%} \)), and we can safely set an accuracy goal of within 20% from satellite measurements after further evaluation and refinement. Remote-sensing products with such a small error will bring confidence on quantitative ocean-color remote sensing and boost interests and applications of professionals and the general public.

[35] Because of these distinctive and unique characteristics, and because the error of measured \( z_{1\%} \) is significantly smaller than that of chlorophyll-a concentration, \( z_{1\%} \) (and/or \( z_{10\%} \)) can be an ideal parameter to validate a system of ocean-color remote sensing. More importantly, \( z_{1\%} \) (or \( z_{10\%} \)) measures water clarity much more rigorously than Secchi depth [Preisendorfer, 1986] and provides much more reliable results. The variations of \( z_{1\%} \) (or \( z_{10\%} \)) (after removing sun angle effects) indicate clearly changes of water quality of an ecosystem. A time series of \( z_{1\%} \) (or \( z_{10\%} \)) of the global oceans from satellite measurements, combined with historical and new measurements, will provide us unprecedented and confident evaluation of the distribution and trend of water clarity of the world oceans. Note that water clarity has profound effects on primary production [Behrenfeld and Falkowski, 1997; Platt and Sathyendranath, 1988] and heat deposition [Kirk, 1988; Lewis et al., 1990; Morel and Antoine, 1994; Zaneveld et al., 1981] in the upper water column.

Appendix A: Updated Quasi-Analytical Algorithm (QAA_v4)

[36] The Quasi-Analytical Algorithm (QAA) was developed by Lee et al. [2002] to derive the absorption and backscattering coefficients by inverting spectral remote-sensing reflectance (\( R_s(\lambda) \)). QAA starts with the calculation of the total absorption coefficient (\( a(\lambda) \)) at a reference wavelength (\( \lambda_0 \)), and then propagate the calculation to other wavelengths. To briefly summarize, this algorithm consists of the following elements:

[37] 1. The ratio of backscattering coefficient (\( b_b \)) to the sum of backscattering and absorption coefficients (\( b_b(a + b_b) \)) at \( \lambda \) is calculated algebraically based on the models of Gordon et al. [1988] and Lee et al. [1999],

\[
\frac{b_b(\lambda)}{a(\lambda) + b_b(\lambda)} = \frac{-0.0895 + \sqrt{0.008 + 0.499R_s(\lambda)}}{0.249}.
\] (A1)

Here \( R_s(\lambda) \) is the spectral remote-sensing reflectance just below the surface and is calculated from \( R_s(\lambda) \) through,

\[
R_s(\lambda) = R_s(\lambda)(0.52 + 1.7R_s(\lambda)).
\] (A2)
2. The spectral \( b_b(\lambda) \) is modeled with the widely used expression [Gordon and Morel, 1983; Smith and Baker, 1981],

\[
b_b(\lambda) = b_{bw}(\lambda) + b_{bp}(\lambda_0) \left( \frac{\lambda_0}{\lambda} \right)^{\eta}.
\]

where \( b_{bw} \) and \( b_{bp} \) are the backscattering coefficients of pure seawater and suspended particles, respectively. Values of \( b_{bw}(\lambda) \) are provided in Morel [1974].

3. When \( \alpha(\lambda_0) \), the ratio of \( b_b(\lambda)/(a + b_b) \) at \( \lambda_0 \), and \( b_{bw}(\lambda_0) \) are known, \( b_{bp}(\lambda_0) \) in equation (A3) can be easily derived with the combination of equations (A1) and (A3). The values of \( b_{bp}(\lambda) \) at other wavelengths are then calculated after the power parameter (\( \eta \)) is estimated from [Lee et al., 2002],

\[
\eta = 2.2 \left(1 - 1.2 \exp \left(-0.9 \frac{r_{rs}(440)}{r_{rs}(555)}\right)\right).
\]

4. Finally, by applying \( b_b(\lambda) \) to the ratio of \( b_b/(a + b_b) \) at \( \lambda \) (equation (A1)), the total absorption coefficient at \( \lambda \), \( \alpha(\lambda) \), is calculated algebraically.

The update of the QAA is related to the calculation of \( \alpha(\lambda_0) \). In this updated version of QAA (v.4), it has two selections for the calculation of \( \alpha(555) \): one for data sets where there are no \( R_{rs} \) measurements at longer wavelengths (640 or 670 nm, for instance); and one for data sets where there are \( R_{rs} \) measurements at those longer wavelengths. And the second round of calculation [Lee et al., 2002] is removed.

For data sets having no \( R_{rs} \) at longer wavelengths, \( \alpha(555) \) is now estimated using the ratio \( \rho = \log(\max(R_{rs}(440, 490, 510))/R_{rs}(555)) \) (for use with the SeaWiFS sensor), and takes the form,

\[
K(555) = 0.0605 + 10^{-1.163 -1.969\rho + 1.239\rho^2 + 0.417\rho^3 - 0.984\rho^4},
\]

\[
\alpha(555) = \frac{0.9K(555)(1 - 6.8R_{rs}(555))}{1 + 15.3R_{rs}(555)}.
\]

The constants in equations (A5) and (A6) are combinations of those from OC4v4 [O’Reilly et al., 2000] and Morel and Maritorena [2001].

For data sets having \( R_{rs} \) at 640 nm (note that \( R_{rs} \) at longer wavelengths are very important for inversion of
waters with high absorption coefficients [IOCCG, 2006; Lee et al., 2005a], $a(555)$ is now estimated as follows,

$$
\chi = \log \left( \frac{R_s(440) + R_s(490)}{R_o(555) + 2 \frac{R_o(640)}{R_o(490)} R_o(640)} \right),
$$

(A7)

$$
a(555) = a_o(555) + 10^{-1.226 - 1.214 z - 0.358 z^2},
$$

(A8)

with constants in equation (A8) obtained by least squares fitting $a(555)$ of the synthetic data set adopted by the IOCCG [IOCCG, 2006].

[45] In the original version of QAA [Lee et al., 2002] a relatively arbitrary boundary was used for the transition between values calculated with 555 nm and 640 nm as reference wavelengths, respectively. In this updated QAA (equations (A7) and (A8)), now seamless transition is achieved for waters from open ocean to coastal region. Available measurements suggest that QAA_v4 is applicable to waters with $a(440)$ as high as $\approx 2.0 \text{ m}^{-1}$. Simply change the corresponding values at 555 nm to that at 550 nm, QAA_v4 (equations (A7) and (A8)) can be applied to data from both SeaWiFS and MODIS without significant effects to the final results.

[46] For the SeaWiFS and MODIS sensor, there is no band to measure $R_s$ at 640 nm, but there are bands to measure $R_s$ at 667 and/or 678 nm. For such cases, a simulated $R_s(640)$ needs to be generated as in Lee et al. [2005a], which is

$$
R_s(640) = 0.01 R_s(555) + 1.4 R_s(667) - 0.0005 R_o(667)/R_o(640),
$$

(A9)

and this simulated $R_s(640)$ is kept greater than 1.2 $R_o(667)$.

[47] Figure A1 compares optical properties derived from $R_s$ by QAA_v4 ($a(555)$ from equation (A8)) with that from water-sample measurements, where data is from the NOMAD_OC_BAM data set [Werdell and Bailey, 2005], but subsampled for valid measurements of both $R_s$ and IOPs (T. Smyth, personal communication). Clearly, excellent agreements were achieved between the two independent determinations. To indicate the applicability of QAA_v4 with satellite data, it was the simulated $R_s(640)$ (equation (A9)) instead of the measured $R_s(640)$ used in the derivation of IOPs (and the $z_{opt}$) in this study.

[48] For a recent measurement made in the south Pacific gyre (M. Lewis, personal communication), consistent results (with $\pm 1\%$ around 150 m) were obtained from measurements of the vertical profiles and the remote sensing reflectance.

[49] Acknowledgments. We thank Tom Peacock, Joseph Rhea, and Jim Mueller for helping with the data collection. Comments and suggestions from Grace Chang and Emmanuel Boss are greatly appreciated. Financial support was provided by NASA through grant NNH06AE55I, and NRL through the 6.1 project “Physical-Biological-Optical Modeling of the Coastal Environment” funded by the Office of Naval Research under Program Element 61153N and 6.2 Project “Tactical UAV hyperspectral imagery for riverine special operations” under Program Element 060245S. Funding was also provided to USF through ONR contract N00014-02-1-0211 and NASA grant NNGO4GL55G.

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